

Decay of solutions for a higher-order parabolic type Kirchhoff equation

Erhan Pişkin¹ and Evrim Akkurt²

¹Dicle University, Department of Mathematics
Diyarbakır, Turkey

²Dicle University, Institute of Natural and Applied Sciences
Department of Mathematics, Diyarbakır, Turkey
Email: episkin@dicle.edu.tr, akkurtevrin@gmail.com

(Received: October 31, 2024 Accepted: Decemeber 13, 2024)

Abstract: In this work, we consider the higher-order parabolic type Kirchhoff equation in a bounded domain. Firstly, we establish the global existence of solutions. Finally, we prove the decay of solutions using the Komornik's lemma.

1 Introduction

1.1 Setting of the problem

In this work, we consider the following initial-boundary value problem for a class of higher-order parabolic-type Kirchhoff equation

$$\left(1 + |z|^{p-2}\right) z_t + M \left(\left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 \right) \mathcal{A}z = |z|^{q-2} z, \quad (x, t) \in \Omega \times (0, T), \quad (1.1)$$

with the initial-boundary conditions

$$z(x, 0) = z_0(x), \quad x \in \Omega,$$

Keywords and phrases: Decay, global existence, Kirchhoff type equation, parabolic-type equation.

2020 AMS Subject Classification: 35A01, 35B40, 35G25.

and

$$\frac{\partial^i z(x, t)}{\partial \nu^i} = 0, \quad i = 0, 1, 2, \dots, m-1, \quad (x, t) \in \partial\Omega \times (0, T),$$

where $\mathcal{A} = (-\Delta)^m$, $m \geq 1$ is a natural number, Ω is a bounded domain in R^n ($n \geq 1$) with smooth boundary $\partial\Omega$, ν is unit outward normal vector on $\partial\Omega$; $p, q > 2$ and

$$M(s) = 1 + s^\gamma, \quad \gamma \geq 1.$$

The problem (1.1) is a generalization of a model introduced by Kirchhoff [7]. More precisely, Kirchhoff proposed a model given by the equation for $f = g = 0$,

$$\rho h z_{tt} + \delta z_t + g(z_t) = \left\{ \rho_0 + \frac{Eh}{2L} \int_0^L (z_x)^2 dx \right\} z_{xx} + f(z), \quad (1.2)$$

for $0 < x < L$, $t \geq 0$, where $z(x, t)$ is the lateral displacement, E the Young modulus, ρ the mass density, h the cross-section area, L the length, ρ_0 the initial axial tension, δ the resistance modulus, and f and g the external forces. Moreover, (1.2) is called a degenerate equation when $\rho_0 = 0$ and nondegenerate one when $\rho_0 > 0$.

1.2 Literature overview

When $m = 1$, (1.1) becomes the following Kirchhoff-type parabolic equation

$$(1 + |z|^{p-2}) z_t - M(\|\nabla z\|^2) \Delta z = |z|^{q-2} z. \quad (1.3)$$

Ouaoua et al. [9] studied the global existence and decay of solutions of the problem (1.3). Later, Khaldi et al. [6] studied by taking variables instead of constants p and q in (1.2).

Pişkin and Ekinçi [14] studied the blow up and growth for the following equation with initial-boundary conditions

$$z_t - \Delta z_t - M(\|\nabla z\|^2) \Delta z + |z|^{q-2} z_t = |z|^{p-2} z.$$

Pişkin and Cömert [12] studied the following nonlinear Kirchhoff-type parabolic equation with logarithmic source term

$$z_t - M(\|\nabla z\|^2) \Delta z - \Delta z_t = |z|^{q-2} z \ln |z|.$$

They established the existence and decay solutions.

Ishige et al. [4] considered the Cauchy problem for nonlinear higher-order heat equation as follows

$$z_t + \mathcal{A}z = |z|^q.$$

They proved the existence of solutions for Cauchy problem.

Xiao and Li [18] considered initial boundary value problem for nonlinear higher-order heat equations of

$$z_t + \mathcal{A}z_t + \mathcal{A}z = f(z).$$

They established the existence of solutions.

Moreover, numerous researchers have studied the existence, decay and blow up of solutions for the higher-order partial differential equations [2, 3, 5, 10, 11, 13, 15, 17, 19, 20]. Motivated by the above studies, in this work, we investigate the global existence and decay solutions for the eq. (1.1). This paper is organized as follows: In Section 2, we introduce some lemmas which will be needed later. In Section 3, we prove the global existence of solutions. In Section 4, we prove the decay of solutions using the Komornik’s lemma.

2 Preliminaries

In this part, we shall give some lemmas, which will be used throughout this work. Let $\|\cdot\|$ and $\|\cdot\|_p$ denote the usual $L^2(\Omega)$ norm and $L^p(\Omega)$ norm, respectively.

Lemma 2.1. (Sobolev-Poincaré inequality) [1, 16] Let p be a number with

$$\begin{cases} 2 \leq p < +\infty, & \text{if } n \leq 2m, \\ 2 \leq p \leq \frac{2n}{n-2m}, & \text{if } n > 2m. \end{cases}$$

Then, there is constant C depending on Ω and p such that

$$\|u\|_p \leq C \left\| \mathcal{A}^{\frac{1}{2}} u \right\|, \forall u \in H_0^m(\Omega).$$

Lemma 2.2. [8]. Let $F : R^+ \rightarrow R^+$ be a non-increasing function and assume that there are two constants $\alpha > 0$ and $C > 0$ such that

$$\int_t^\infty F^{\alpha+1}(s) ds \leq CF^\alpha(0) F(t), \forall t \in R^+.$$

Then,

$$F(t) \leq \eta(C + \alpha t)^{\frac{-1}{\alpha}}, \quad \forall t \geq C,$$

where η is a positive constant.

3 Global existence

In this part, we prove the global existence of solutions.

The energy function $E(t)$ and the Nehari's functional $I(t)$ can be defined respectively by

$$E(t) = E(z) = \frac{1}{2} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \frac{1}{2(\gamma+1)} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} - \frac{1}{q} \|z\|_q^q \quad (3.1)$$

and

$$I(t) = I(z) = \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} - \|z\|_q^q. \quad (3.2)$$

Lemma 3.1. *Suppose that z be a solution of (1.1). Then, the energy function $E(t)$ is a nonincreasing function for $t \in [0, T]$ and*

$$E'(t) = -\|z_t\|^2 - \int_{\Omega} |z|^{p-2} |z_t|^2 dx \leq 0, \quad (3.3)$$

and

$$E(t) \leq E(0).$$

Proof. Multiplying the equation in (1.1) by z_t and integrating over the domain Ω , to obtain

$$\frac{d}{dt} \left[\frac{1}{2} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \frac{1}{2(\gamma+1)} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} - \frac{1}{q} \|z\|_q^q \right] = -\|z_t\|^2 - \int_{\Omega} |z|^{p-2} |z_t|^2 dx.$$

Then,

$$E'(t) = -\|z_t\|^2 - \int_{\Omega} |z|^{p-2} |z_t|^2 dx \leq 0.$$

Integrating (3.3) over $(0, t)$, we get

$$E(t) \leq E(0).$$

□

Lemma 3.2. *Let z be a solution of (1.1) and $q > 2(\gamma + 1)$,*

$$I(0) > 0,$$

and

$$\beta_1 + \beta_2 < 1, \quad (3.4)$$

here

$$\beta_1 = \alpha c_*^q \left(\frac{2q}{q-2} E(0) \right)^{\frac{q-2}{q}} \quad \text{and} \quad \beta_2 = (1-\alpha) c_*^q \left(\frac{2(\gamma+1)q}{q-2(\gamma+1)} E(0) \right)^{\frac{q-2(\gamma+1)}{2(\gamma+1)}},$$

with $0 < \alpha < 1$, and c_* is the best embedding constant of $H_0^m(\Omega) \hookrightarrow L^q(\Omega)$, then

$$I(t) > 0,$$

for all $t \in [0, T]$.

Proof. By continuity, there exist T_* such that

$$I(t) \geq 0, \quad \text{for all } t \in [0, T_*]. \quad (3.5)$$

By direct computation, we obtain for all $t \in [0, T]$,

$$\begin{aligned} E(t) &= E(z) = \frac{1}{2} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \frac{1}{2(\gamma+1)} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} - \frac{1}{q} \|z\|_q^q \\ &= \frac{1}{2} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \frac{1}{2(\gamma+1)} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} \\ &\quad - \frac{1}{q} \left(\left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} - I(t) \right) \\ &\geq \frac{q-2}{2q} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \frac{q-2(\gamma+1)}{2(\gamma+1)q} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} + \frac{1}{q} I(t). \end{aligned}$$

By using (3.5), we have

$$\frac{q-2}{2q} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \frac{q-2(\gamma+1)}{2(\gamma+1)q} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} \leq E(t), \quad \forall t \in [0, T_*]. \quad (3.6)$$

By the definition of $E(t)$, we obtain

$$\begin{aligned} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 &\leq \frac{2q}{q-2} E(t) \\ &\leq \frac{2q}{q-2} E(0) \end{aligned} \quad (3.7)$$

and

$$\begin{aligned} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} &\leq \frac{2q(\gamma+1)}{q-2(\gamma+1)} E(t) \\ &\leq \frac{2q(\gamma+1)}{q-2(\gamma+1)} E(0). \end{aligned} \quad (3.8)$$

On the other hand, we get

$$\|z\|_q^q = \alpha \|z\|_q^q + (1-\alpha) \|z\|_q^q.$$

By the embedding of $H_0^m(\Omega) \hookrightarrow L^q(\Omega)$, we have

$$\begin{aligned} \|z\|_q^q &\leq \alpha c_*^q \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^q + (1-\alpha) \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^q \\ &= \alpha c_*^q \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{q-2} \cdot \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 \\ &\quad + (1-\alpha) \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{q-2(\gamma+1)} \cdot \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)}. \end{aligned}$$

From (3.7) and (3.8), we have

$$\begin{aligned} \int_{\Omega} |z|^q dx &\leq \alpha c_*^q \left(\frac{2q}{q-2} E(0) \right)^{\frac{q-2}{2}} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 \\ &\quad + (1-\alpha) c_*^q \left(\frac{2q(\gamma+1)}{q-2(\gamma+1)} E(0) \right)^{\frac{q-2(\gamma+1)}{2(\gamma+1)}} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)}, \forall t \in [0, T_*], \end{aligned}$$

which implies that

$$\int_{\Omega} |z|^q dx \leq \beta_1 \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \beta_2 \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)}, \forall t \in [0, T_*]. \quad (3.9)$$

Since $\beta_1 + \beta_2 < 1$, so

$$\int_{\Omega} |z|^q dx \leq \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)}, \quad \forall t \in [0, T_*].$$

This implies that

$$I(t) > 0, \quad \forall t \in [0, T_*].$$

By repeating the above procedure, we can extend T_* to T . \square

Theorem 3.1. *Under the assumptions of above lemma, the local solution of (1.1) is global.*

Proof. It is sufficient to show that $\left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)}$ is bounded independently of t . By using (3.1) and (3.2), we have

$$\begin{aligned} E(t) &= \frac{1}{2} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \frac{1}{2(\gamma+1)} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} - \frac{1}{q} \|z\|_q^q \\ &= \frac{q-2}{2q} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \frac{q-2(\gamma+1)}{2(\gamma+1)q} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} + \frac{1}{q} I(t). \end{aligned}$$

Since $I(t) \geq 0$, which implies that

$$\left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} \leq CE(t),$$

so

$$\left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 + \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2(\gamma+1)} \leq CE(0),$$

where $C = \max \left\{ \frac{2q}{q-2}, \frac{2(\gamma+1)q}{q-2(\gamma+1)} \right\}$. \square

4 Decay of solutions

In this part, we prove the decay of solutions by using the Komornik's lemma.

Lemma 4.1. *Assume that the assumptions of Lemma 3.1 and $p > 2$, hold. Then, there exist a positive constant c such that*

$$\|z\|_p^p \leq cE(t). \quad (4.1)$$

Proof. By direct computation, we get

$$\begin{aligned}
\|z\|_p^p &\leq c_*^p \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^p \\
&\leq c_*^p \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{p-2} \cdot \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^2 \\
&\leq c_*^p \left(\frac{2q}{q-2} E(0) \right)^{\frac{p-2}{2}} \cdot \frac{2q}{q-2} E(t) \\
&\leq cE(t).
\end{aligned}$$

□

Theorem 4.1. *Suppose that the assumptions of Lemma 3.1. Then, there exists constants $C > 0$, such that*

$$E(t) \leq E(0) \left(\frac{C + qt}{C + qC} \right)^{\frac{-1}{q}}, \text{ for all } t \geq C.$$

Proof. Multiplying the equation of (1.1) by $z(t) E^q(t)$, ($q > 0$) and integrating over $\Omega \times (S, T)$, we get

$$\begin{aligned}
&\int_S^T \int_{\Omega} E^q(t) \left[z z_t - z \left(M \left(\int_{\Omega} \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 dx \right) \mathcal{A} z + |z_t|^{p-2} z \right) \right] dx dt \\
&= \int_S^T E^q(t) \int_{\Omega} |z|^q dx dt.
\end{aligned}$$

It follows that

$$\begin{aligned}
&\int_S^T \int_{\Omega} E^q(t) \left[z z_t + \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 + \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2\gamma} \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 + z |z|^{p-2} z_t \right] dx dt \\
&= \int_S^T E^q(t) \int_{\Omega} |z|^q dx dt.
\end{aligned}$$

We add and subtract the term

$$\int_S^T E^q(t) \int_{\Omega} \left[\beta_1 \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 + \beta_2 \left\| \mathcal{A}^{\frac{1}{2}} z \right\|_2^{2\gamma} \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 \right] dx dt$$

and use (3.9) to obtain

$$\begin{aligned}
 & (1 - \beta_1) \int_S^T E^q(t) \int_{\Omega} \left[\left| \mathcal{A}^{\frac{1}{2}} z \right|^2 dx dt \right] \\
 & + (1 - \beta_2) \int_S^T E^q(t) \int_{\Omega} \left[\left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2\gamma} \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 \right] dx dt \\
 & + \int_S^T E^q(t) \int_{\Omega} [zz_t] dx dt \\
 & + \int_S^T E^q(t) \int_{\Omega} zz_t |z|^{p-2} dx dt \\
 & = - \int_S^T E^q(t) \int_{\Omega} \left[\beta_1 \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 + \beta_2 \left\| \mathcal{A}^{\frac{1}{2}} z \right\|_2^{2\gamma} \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 - |z| \right] dx dt \leq 0. \quad (4.2)
 \end{aligned}$$

It is clear that

$$\begin{aligned}
 & \xi \int_S^T E^q(t) \int_{\Omega} \left[\frac{1}{2} \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 + \frac{1}{2(\gamma+1)} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2\gamma} \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 - \frac{|z|^q}{q} \right] dx dt \\
 & \leq (1 - \beta_1) \int_S^T E^q(t) \int_{\Omega} \left[\frac{1}{2} \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 \right] dx dt \\
 & + (1 - \beta_2) \int_S^T E^q(t) \int_{\Omega} \left[\frac{1}{2(\gamma+1)} \left\| \mathcal{A}^{\frac{1}{2}} z \right\|^{2\gamma} \left| \mathcal{A}^{\frac{1}{2}} z \right|^2 \right] dx dt \quad (4.3)
 \end{aligned}$$

where $\xi = \min \{1 - \beta_1, 1 - \beta_2\}$. By use of (4.2), (4.3) and definition of $E(t)$, we have

$$\begin{aligned}
 \xi \int_S^T E^{q+1}(t) dt & \leq - \int_S^T E^q(t) \int_{\Omega} zz_t dx dt \\
 & \quad - \int_S^T E^q(t) \int_{\Omega} zz_t |z|^{p-2} dx dt. \quad (4.4)
 \end{aligned}$$

We estimate the terms in the right-hand side of (4.4), for the first term, we use the following Young inequality, we have

$$XY \leq \frac{\varepsilon}{\lambda_1} X^{\lambda_1} + \frac{1}{\lambda_2 \varepsilon^{\lambda_1}} Y^{\lambda_2}, \quad X, Y \geq 0, \quad \varepsilon > 0$$

and $\frac{1}{\lambda_1} + \frac{1}{\lambda_2} = 1$,

$$\begin{aligned} & - \int_S^T E^q(t) \int_{\Omega} z z_t dx dt \\ & \leq \int_S^T E^q(t) \int_{\Omega} \left(\varepsilon c |z|^2 + c_{\varepsilon} |z_t|^2 \right) dx dt. \end{aligned} \quad (4.5)$$

We use again the above Young inequality, we get

$$\begin{aligned} & - \int_S^T E^q(t) \int_{\Omega} z z_t |z|^{p-2} dx dt \\ & = - \int_S^T E^q(t) \int_{\Omega} |z|^{\frac{p-2}{2}} z_t |z|^{\frac{p-2}{2}} z dx dt \\ & \leq \int_S^T E^q(t) \int_{\Omega} \left(\varepsilon c |z|^p + c_{\varepsilon} |z|^{p-2} z_t^2 \right) dx dt. \end{aligned} \quad (4.6)$$

By (4.5) and (4.6), inequality (4.4), becomes

$$\begin{aligned} \xi \int_S^T E^{q+1}(t) dt & \leq \int_S^T E^q(t) \int_{\Omega} \left(\varepsilon c |z|^2 + c_{\varepsilon} |z|^2 \right) dx dt \\ & \quad + \int_S^T E^q(t) \int_{\Omega} \left(\varepsilon c |z|^p + c_{\varepsilon} |z|^{p-2} z_t^2 \right) dx dt \\ & \leq \varepsilon c \int_S^T E^q(t) \int_{\Omega} \left[|z|^2 + |z|^p \right] dx dt \\ & \quad + c_{\varepsilon} \int_S^T E^q(t) \int_{\Omega} \left(|z_t|^2 + |z|^{p-2} z_t^2 \right) dx dt. \end{aligned} \quad (4.7)$$

By (3.10) and definition of $E'(t)$, we have

$$\xi \int_S^T E^{q+1}(t) dt \leq \varepsilon c \int_S^T E^{q+1}(t) dt + c_{\varepsilon} \int_S^T E^q(t) \int_{\Omega} (-E'(t)) dx dt. \quad (4.8)$$

This implies

$$\begin{aligned} \xi \int_S^T E^{q+1}(t) dt &\leq \varepsilon c \int_S^T E^{q+1}(t) dt + c_\varepsilon [E^{q+1}(S) - E^{q+1}(T)] \\ &\leq \varepsilon c \int_S^T E^{q+1}(t) dt + c_\varepsilon E^q(0) E(S). \end{aligned} \quad (4.9)$$

Choosing ε small enough such that, we arrive at

$$\int_S^T E^{q+1}(t) dt \leq c E^q(0) E(S).$$

By taking T goes to ∞ , we have

$$\int_S^\infty E^{q+1}(t) dt \leq c E^q(0) E(S).$$

Komornik's integral inequality yields the desired result. \square

Acknowledgment:

The author is thankful to the editor(s) and referee(s) for their valuable suggestions to improve the manuscript.

References

- [1] R. A. Adams and J. J. F. Fournier, *Sobolev Spaces*, Academic Press, New York, 2003.
- [2] T. Cömert and E. Pişkin, *Blow up at Infinity of Weak Solutions for a Higher-Order Parabolic Equation with Logarithmic Nonlinearity*, J. Uni. Math., **4**(2)(2021), 118-127.
- [3] T. Cömert and E. Pişkin, *Global Existence and Exponential Decay of Solutions for Higher-Order Parabolic Equation with Logarithmic Nonlinearity*, Miskolc Math. Notes, **23**(2)(2022), 595-605.
- [4] K. Ishige, T. Kawakami and S. Okabe, *Existence of solutions for a higher-order semilinear parabolic equation with singular initial data*, Ann. Inst. H. Poincaré Anal. non linéaire, **37**(5)(2020), 1185-1209.

- [5] N. İrkıl and E. Pişkin, *Global existence and decay of solutions for a higher-order Kirchhoff-type systems with logarithmic nonlinearities*, Quaestiones Math., **45**(4)(2022), 523-546.
- [6] A. Khaldi, A. Ouaoua and M. Maouni, *Global existence and stability of solution for a nonlinear Kirchhoff type reaction-diffusion equation with variable exponents*, Math. Bohem., **147**(4)(2022), 471-484.
- [7] G. Kirchhoff, *Vorlesungen uber mechanik*, BG Teubner, 1883.
- [8] V. Komornik, *Exact controllability and stabilization the multiplier method*, Paris: Masson-John Wiley, 1994.
- [9] A. Ouaoua, A. Khaldi and M. Maouni, *Stabilization of solutions for a Kirchhoff type reaction-diffusion equation*, Canad. J. Appl. Math., **2**(2)(2020), 71-80.
- [10] E. Pişkin, *Blow up Solutions for a Class of Nonlinear Higher-Order Wave Equation with Variable Exponents*, Sigma J. Eng. Nat. Sci., **10**(2)(2019), 149-156.
- [11] E. Pişkin and R. Aksoy, *Higher Order Logarithmic Klein-Gordon Equation: Global Existence, Decay and Nonexistence*, Acta Univ. Apulensis Math., **70**(2022), 47-63.
- [12] E. Pişkin and T. Cömert, *Existence and decay of solutions for a parabolic type Kirchhoff equation with logarithmic nonlinearity*, Adv. Stud. Euro-Tbil. Math. J., **15**(1)(2022), 111-128.
- [13] E. Pişkin and A. Demirhan, *On potential wells for a nonlinear higher-order hyperbolic-type equation*, Int. J. Open Probl Comput. Sci. Math., **17**(2)(2024), 24-37.
- [14] E. Pişkin and F. Ekinçi, *Blow up and growth of solutions for a parabolic type Kirchhoff equation with multiple nonlinearities*, Konuralp J. Math., **8**(1)(2020), 216-222.
- [15] E. Pişkin and N. İrkıl, *Existence and Decay of Solutions for a Higher-Order Viscoelastic Wave Equation with Logarithmic Nonlinearity*, Commun. Fac. Sci. Univ. Ank. Ser. A1. Math. Stat., **70**(1)(2021), 300-319.
- [16] E. Pişkin and B. Okutmuşur, *An Introduction to Sobolev Spaces*, Bentham Science, 2021.

- [17] M. Shahrouzi, *Asymptotic behavior of solutions for a nonlinear viscoelastic higher-order $p(x)$ -Laplacian equation with variable-exponent logarithmic source term*, Bol.Soc. mat. mex., **29**(2023), 1-20.
- [18] L. Xiao and M. Li, *Initial boundary value problem for a class of higher-order n -dimensional nonlinear pseudo-parabolic equations*, Bound. Value Probl., 2021, (2021), 1-24.
- [19] L. M. Xiao, C. Luo and J. Liu, *Global existence of weak solutions to a class of higher-order nonlinear evolution equations*, Electron. Res. Arch., **32**(9)(2024), 5357-5376.
- [20] H. Yükksekaya, E. Pişkin, S. M. Boulaaras and B. B. Cherif, *Existence, Decay and Blow-up of Solutions for a Higher-Order Kirchoff-Type Equation with Delay Term*, J. Funct. Spaces, (2021), 1-11.